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# Vibrational assignments and derived potential energy distributions for tri- and difluoromethyl ketene by density functional calculations 

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#### Abstract

The structures of 3,3,3-trifluoromethyl ketene and 3,3 -difluoromethyl ketene were studied by utilizing ab initio calculations with the $6-311++\mathrm{G}^{* *}$ basis set at the (B3LYP) Density Functional level. Full optimization was performed for both molecules in their ground and transition states. Energy optimization of the systems under investigation shows that trifluoromethyl ketene exists only in the cis conformation (fluorine atom eclipses the ketene group). Difluoromethyl ketene was predicted to have two stable conformations: the cis (hydrogen atom eclipses the ketene group) and the gauche (fluorine atom eclipses the ketene group) form. The conformational stability of the molecules was found to be governed mainly by electrostatic and molecular orbital interactions. The vibrational frequencies were computed and complete assignments were provided on the basis of normal coordinate calculations and comparison with similar molecules. The potential energy distributions (PED) among symmetry coordinates were derived for the stable conformations of the two molecules.


Keywords Vibrational spectra and assignments . Rotational barriers • Trifluoromethyl ketene Difluoromethyl ketene

## Introduction

The interesting electronic structure and chemical reactivity of ketenes have been demonstrated in the literature. [1, 2, 3, 4, 5, 6, 7, 8, 9] Many simple substituted ketenic molecules have been isolated experimentally, characterized, and used to produce other ketenic and organic systems. $[10,11,12,13,14,15,16,17,18,19]$ The importance of ketenes as well as aldehydes [20, 21, 22, 23, 24, $25,26,27,28,29,30,31,32,33,34]$ and isocyanates

[^0]$[35,36,37,38,39,40,41,42,43,44,45,46,47,48,49$, 50,51] in synthetic chemistry has attracted attention in order to study their properties and structural stabilities for many years.

When conjugation is the predominant force, as in acrolein, [20,23] formyl ketene [9] and formyl isocyanate, [51] the molecules are found to exist as a mixture of the planar cis and trans conformations. In these compounds, it is clear that the $\pi-\pi$ interaction between the $\mathrm{C}=\mathrm{O}$ bond and vinyl $-\mathrm{C}=\mathrm{C}-$, the ketene $-\mathrm{C}=\mathrm{C}=\mathrm{O}$ or the isocyanate $-\mathrm{N}=\mathrm{C}=\mathrm{O}$ groups greatly stabilizes the planar forms with a relatively high rotational barrier. In comparison, the $\mathrm{C}-\mathrm{N}$ barrier in the isocyanates is much lower than the $\mathrm{C}-\mathrm{C}$ barrier in the vinyl derivatives and ketenes. This significant difference between the two rotational barriers is due to the difference in the strength of the partial $\pi$ character of the $\mathrm{C}-\mathrm{N}$ bond as compared to that of the $\mathrm{C}-\mathrm{C}$ bond. This difference is the result of the electronegativity change on going from the $\mathrm{C}-\mathrm{C}$ to the $\mathrm{C}-\mathrm{N}$ bond.

In the case of halomethyl isocyanates, the molecules were found to exist in only one conformation with minimum interaction between the halogen atom and the lone pair. [49, 50] For example, for chloromethyl isocyanate the trans conformation (chlorine atom eclipses the nitrogen lone pair) was found to be the stable conformation. [49] Difluoromethyl isocyanate was also predicted to exist in the trans conformation in which the hydrogen atom eclipses the lone pair. [50] In haloacetaldehyde systems, [24,25] it was found that the planar s-trans conformation (the halogen atom eclipses the aldehydic hydrogen) is the form with the lowest energy. The steric effect (the repulsive force between the halogen atom and the carbonyl group) would destabilize the cis form of the molecule.

3,3,3-Trifluoromethyl ketene has been reported to be an important intermediate. [52, 53] The rate of hydration of the molecule was observed experimentally using UV spectroscopy. [53] Moreover, the reactivity of trifluoromethyl ketene was studied using molecular orbital methods, and the substituent effect of the $\mathrm{CF}_{3}$ group in the
system was compared to a series of other substituents theoretically. [3] As a continuation of these studies, we investigated the structure of 3,3,3-trifluoromethyl ketene and 3,3-difluoromethyl ketene to understand the nature of the forces that control the conformational behavior of such important molecular systems. We carried out normal coordinate calculations by utilizing ab initio density functional calculations and derived the potential energy distributions among symmetry coordinates. The vibrational assignments of the normal modes were made for both molecules. The results of the work are presented herein.

## Ab initio calculations

The GAUSSIAN 98 program [54] running on an IBM RS/6000 43P model 260 workstation, was used to carry out the LCAO-MO-SCF DFT-B3LYP calculations. The extended $6-311++G^{* *}$ basis set was employed to optimize the structures and predict the energies and dipole moments of trifluoromethyl ketene and difluoromethyl ketene in their stable conformations. The calculations were extended to the Density Functional B3LYP level. From full energy optimization the cis conformation of trifluoromethyl ketene was found to be the energy minimum for the system, whereas both the cis and the gauche conformations were the stable conformers in difluoromethyl ketene (Fig. 1). The calculated structural parameters in the molecules are shown in Tables 1 and 2. Moreover, the optimized geometrical parameters of trifluoromethyl ketene were compared to the corresponding values obtained from experimental data for trifluoroacetaldehyde [55] as in Table 1. The optimized structural parameters were used to compute vibrational frequencies at the DFT-B3LYP/6-311++G** level. Normal coordinate calculations were then carried out to derive the potential energy distributions for the molecules in their stable conformations.

## Asymmetric torsional potential function

The potential scan for the internal rotation about the $\mathrm{C}-\mathrm{C}$ single bond was obtained by allowing the OCCC dihedral angle $(\phi)$ to vary from $0^{\circ}$ (cis position) to $180^{\circ}$ (trans position). Full geometry optimizations at each of the fixed dihedral angles $(\phi)$ of $15^{\circ}, 30^{\circ}, 45^{\circ}, 60^{\circ}, 75^{\circ}$, $90^{\circ}, 105^{\circ}, 120^{\circ}, 135^{\circ}, 150^{\circ}$, and $165^{\circ}$ were carried out at DFT-B3LYP/6-311++G** level in both systems. The barrier to interconversion and the relative energies in difluoromethyl ketene were calculated and are shown in Table 3. The torsional potential was represented as a Fourier cosine series in the dihedral angle ( $\phi$ ): $V(\phi)=V_{0}+\Sigma\left(V_{n} / 2\right)[1-\cos (n \phi)]$, where $V_{0}$ is the relative energy at the cis conformation, and the potential coefficients from $V_{1}$ to $V_{6}$ are considered adequate to describe the potential function. The results of the energy optimizations were used to calculate the six coefficients by least-squares fitting for the systems (Table 4). The poten-





Fig. 1 Atom numbering for trifluoromethyl ketene (upper) and difluoromethyl ketene (lower) in the cis (right) and the gauche (left) conformations

Table 1 Calculated structural parameters, total dipole moment, and rotational constants for the cis conformation of trifluoromethyl ketene


Dipole moment (Debye)
$\mu_{\mathrm{t}} \quad 2.1$
Rotational constants (MHz)

| $A$ | 5428 |
| :--- | :--- |
| $B$ | 1681 |
| $C$ | 1657 |

${ }^{\text {a }}$ Calculated total energy, zero-point correction, and the sum of total and zero point energies (in Hartrees) of the cis conformer of trifluoromethyl ketene at DFT-B3LYP/6-311++G** level are $-489.79721,0.03794$, and -489.75927 , respectively
${ }^{\mathrm{b}}$ Data was obtained for trifluoromethylacetaldehyde (fluoral), see ref [55]
tial functions of trifluoromethyl ketene and difluoromethyl ketene are shown in Fig. 2.

Vibrational frequencies and normal coordinate analyses
Both trifluoromethyl ketene and difluoromethyl ketene in their cis conformations have $\mathrm{C}_{\mathrm{s}}$ symmetry. The 18 vibrational modes span the irreducible representations: 12 $A^{\prime}$ and $6 A^{\prime \prime}$. The A' modes should be polarized while the A" modes are depolarized in the Raman spectra of the liquid. However, the gauche conformation in difluoromethyl ketene has $\mathrm{C}_{1}$ symmetry. All the eighteen vibrational modes in difluoromethyl ketene belong to the A

Table 2 Calculated structural parameters, total dipole moment, and rotational constants for the cis and trans conformations of difluoromethyl ketene

| Parameter | B3LYP/6-311++G** |  |
| :---: | :---: | :---: |
|  | Cis | Gauche |
| Bond length ( A ) |  |  |
| $r\left(\mathrm{C}_{1}-\mathrm{C}_{2}\right)$ | 1.484 | 1.489 |
| $r\left(\mathrm{C}_{2}-\mathrm{C}_{3}\right)$ | 1.314 | 1.316 |
| $r\left(\mathrm{C}_{3}-\mathrm{O}_{4}\right)$ | 1.157 | 1.156 |
| $r\left(\mathrm{C}_{2}-\mathrm{H}_{5}\right)$ | 1.083 | 1.082 |
| $r\left(\mathrm{C}_{1}-\mathrm{H}_{6}\right)$ | 1.091 | 1.092 |
| $r\left(\mathrm{C}_{1}-\mathrm{F}_{7}\right)$ | 1.377 | 1.380 |
| $r\left(\mathrm{C}_{1}-\mathrm{F}_{8}\right)$ | 1.377 | 1.372 |
| Bond angle (deg) |  |  |
| $\left(\mathrm{C}_{1} \mathrm{C}_{2} \mathrm{C}_{3}\right)$ | 121.859 | 122.084 |
| $\left(\mathrm{C}_{2} \mathrm{C}_{3} \mathrm{O}_{4}\right)$ | 179.996 | 180.802 |
| $\left(\mathrm{C}_{1} \mathrm{C}_{2} \mathrm{H}_{5}\right.$ ) | 119.425 | 119.629 |
| $\left(\mathrm{C}_{2} \mathrm{C}_{1} \mathrm{H}_{6}\right.$ ) | 114.020 | 113.460 |
| $\left(\mathrm{C}_{2} \mathrm{C}_{1} \mathrm{~F}_{7}\right)$ | 110.651 | 111.223 |
| $\left(\mathrm{C}_{2} \mathrm{C}_{1} \mathrm{~F}_{8}\right)$ | 110.651 | 110.579 |
| $\left(\mathrm{C}_{3} \mathrm{C}_{2} \mathrm{C}_{1} \mathrm{H}_{6}\right)$ | 0.0 | 131.508 |
| $\left(\mathrm{C}_{3} \mathrm{C}_{2} \mathrm{C}_{1} \mathrm{~F}_{7}\right)$ | 121.707 | 119.954 |
| $\left(\mathrm{C}_{3} \mathrm{C}_{2} \mathrm{C}_{1} \mathrm{~F}_{8}\right)$ | -121.717 | -121.680 |
| Dipole moment (Debye) |  |  |
| $\mu_{t}$ | 1.4 | 2.8 |
| Rotational constants (MHz) |  |  |
| A | 8815 | 7928 |
| B | 1906 | 2259 |
| C | 1690 | 1962 |

representation and are polarized in the Raman spectrum of the liquid.

Normal coordinate analyses were carried out for the stable conformers of the molecules in order to provide a complete assignment of the fundamental vibrational frequencies. A computer program was written for this purpose by following Wilson's method. [56] The cartesian coordinates for the stable conformers together with the normal modes (in cartesian coordinates) and the frequencies from the GAUSSIAN 98 output were used as an input in the program. A complete set of internal coordinates (Tables 5 and 6) was used to form symmetry coordinates (Tables 7 and 8 ) in our molecular systems.

The normal modes were next transformed to massweighted cartesian coordinates, which were then used to calculate the force constant matrix. This was diagonalized and its eigenvectors and eigenvalues were used in the further calculations. Following this step the force constant matrix was transformed to internal coordinates. To ensure correctness, this transformation was checked numerically in both directions. At this point the force constant matrix in internal coordinates could be scaled if desired, back-transformed to mass weighted cartesians and diagonalized again to get scaled frequencies and normal modes. The matrix was finally transformed to symmetry coordinates where again all possible numerical checks were performed.

In the next step the normal modes were also transformed to symmetry coordinates. Finally, the potential energy distribution (PED) for each normal mode among

Table 4 Calculated potential constants ( $\mathrm{kcal} \mathrm{mol}^{-1}$ ) for the asymmetric torsion in trifluoromethyl ketene and difluoromethyl ketene

| Potential constants | B3LYP/6-311++G** <br>  <br> Trifluoromethyl <br> ketene $^{\text {a }}$Difluoromethyl <br> ketene |  |
| :--- | :--- | :---: |
| $V_{1}$ |  | 0.115 |
| $V_{2}$ | 1.315 | 1.077 |
| $V_{3}$ |  | 1.289 |
| $V_{4}$ | 0.012 |  |
| $V_{5}$ | -0.021 |  |
| $V_{6}$ |  | -0.022 |

${ }^{\text {a }}$ Because all potential constants, except $V_{3}$, related to the trifluoroderivative are of a very low value, they were omitted

Table 3 Computed total, zeropoint corrections (Hartrees), relative energies, and rotational barriers ( $\mathrm{kcal} \mathrm{mol}^{-1}$ ) in difluoromethyl ketene
${ }^{\text {a }}$ The CCCO dihedral angle of the gauche conformer was calculated to be $131.51^{\circ}$
${ }^{\mathrm{b}}$ The CCCO dihedral angle of the transition state was calculated to be $67.57^{\circ}$

|  | B3LYP/6-311++G** |  |  |
| :--- | ---: | ---: | ---: |
|  | Cis | Gauche $^{\mathrm{a}}$ | TS ${ }^{\mathrm{b}}$ |
| Total energy | -390.51751 | -390.51622 | -390.51402 |
| Relative energy | 0.80949 |  |  |
| cis-gauche barrier | 2.19001 |  |  |
| gauche-cis barrier | 1.38052 |  |  |
| Zero-point correction | 0.04634 | 0.04629 | 0.04599 |
| Sum of total and zero-point energies | -390.47118 | -390.46993 | -390.46803 |
| Corrected relative energy | 0.78439 |  |  |
| Corrected cis-gauche barrier | 1.97666 |  |  |
| Corrected gauche-cis barrier | 1.19227 |  |  |

Fig. 2 Potential curve for the internal rotation in trifluoromethyl ketene (solid line; note, that the "gauche" minimum at $120^{\circ}$ is identical to the cis form) and difluoromethyl ketene (dashed line) as determined by ab initio calculations at B3LYP/6-311++G** level


Table 5 Internal coordinate definitions (see Fig. 1 for atom denotation) for trifluoromethyl ketene

| No. | Coordinate |  | Definition |
| :---: | :--- | :--- | :--- |
| 1 | $\mathrm{C}_{1}-\mathrm{C}_{2}$ | stretch | R |
| 2 | $\mathrm{C}_{2}-\mathrm{C}_{3}$ | stretch | P |
| 3 | $\mathrm{C}_{3}-\mathrm{O}_{4}$ | stretch | T |
| 4 | $\mathrm{C}_{2}-\mathrm{H}_{5}$ | stretch | $\mathrm{D}^{2}$ |
| 5 | $\mathrm{C}_{1}-\mathrm{F}_{6}$ | stretch | $\mathrm{Q}_{1}$ |
| 6 | $\mathrm{C}_{1}-\mathrm{F}_{7}$ | stretch | $\mathrm{Q}_{2}$ |
| 7 | $\mathrm{C}_{1}-\mathrm{F}_{8}$ | stretch | $\mathrm{Q}_{3}$ |
| 8 | $\mathrm{C}_{1} \mathrm{C}_{2} \mathrm{C}_{3}$ | bend | $\gamma$ |
| 9 | $\mathrm{C}_{2} \mathrm{C}_{3} \mathrm{O}_{4}$ | bend | $\sigma$ |
| 10 | $\mathrm{C}_{1} \mathrm{C}_{2} \mathrm{H}_{5}$ | bend | $\rho$ |
| 11 | $\mathrm{C}_{3} \mathrm{C}_{2} \mathrm{H}_{5}$ | bend | $\mathrm{H}^{2}$ |
| 12 | $\mathrm{C}_{2} \mathrm{C}_{1} \mathrm{~F}_{6}$ | bend | $\delta$ |
| 13 | $\mathrm{C}_{2} \mathrm{C}_{1} \mathrm{~F}_{7}$ | bend | $\alpha_{1}$ |
| 14 | $\mathrm{C}_{2} \mathrm{C}_{1} \mathrm{~F}_{8}$ | bend | $\alpha_{2}$ |
| 15 | $\mathrm{~F}_{6} \mathrm{C}_{1} \mathrm{~F}_{7}$ | bend | $\alpha_{3}$ |
| 16 | $\mathrm{~F}_{6} \mathrm{C}_{1} \mathrm{~F}_{8}$ | bend | $\beta_{1}$ |
| 17 | $\mathrm{~F}_{7} \mathrm{C}_{1} \mathrm{~F}_{8}$ | bend | $\beta_{2}$ |
| 18 | $\mathrm{C}_{2}=\mathrm{C}_{3}=\mathrm{O}_{4}$ | wag | $\beta_{3}$ |
| 19 | $\mathrm{H}_{5} \mathrm{C}_{1} \mathrm{C}_{2} \mathrm{C}_{3}$ | wag | $\chi$ |
| 20 | $\mathrm{~F}_{6} \mathrm{C}_{1} \mathrm{C}_{2} \mathrm{C}_{3}$ | (CF $\left.\mathrm{F}_{3}\right)$ torsion | $\pi$ |
|  | $\mathrm{F}_{7} \mathrm{C}_{1} \mathrm{C}_{2} \mathrm{C}_{3}$ |  | $\tau$ |
|  | $\mathrm{~F}_{8} \mathrm{C}_{1} \mathrm{C}_{2} \mathrm{C}_{3}$ |  |  |

the symmetry coordinates was calculated and is given in Tables 9,10 and 11. A complete vibrational assignment of the fundamentals was proposed. The assignments were made based on calculated PED, infrared band intensities, Raman line activities and depolarization ratios. The data of the vibrational assignments are also given in Tables 9, 10 and 11.

## Calculation of vibrational spectra

To calculate the Raman spectra we used the frequencies $v_{j}$, the scattering activities $S_{j}$ and the depolarization ratios $\rho_{j}$ as calculated on DFT-B3LYP level $\left(6-311++G^{* *}\right.$ basis set for all). Then, the Raman cross-sections

Table 6 Internal coordinate definitions (see Fig. 1 for atom denotation) for difluoromethyl ketene

| No. | Coordinate |  | Definition |
| :---: | :---: | :---: | :---: |
| 1 | $\mathrm{C}_{1}-\mathrm{C}_{2}$ | stretch | R |
| 2 | $\mathrm{C}_{2}-\mathrm{C}_{3}$ | stretch | P |
| 3 | $\mathrm{C}_{3}-\mathrm{O}_{4}$ | stretch | T |
| 4 | $\mathrm{C}_{2}-\mathrm{H}_{5}$ | stretch | D |
| 5 | $\mathrm{C}_{1}-\mathrm{H}_{6}$ | stretch | $\mathrm{Q}_{1}$ |
| 6 | $\mathrm{C}_{1}-\mathrm{F}_{7}$ | stretch | $\mathrm{Q}_{2}$ |
| 7 | $\mathrm{C}_{1}-\mathrm{F}_{8}$ | stretch | $\mathrm{Q}_{3}$ |
| 8 | $\mathrm{C}_{1} \mathrm{C}_{2} \mathrm{C}_{3}$ | bend | $\gamma$ |
| 9 | $\mathrm{C}_{2} \mathrm{C}_{3} \mathrm{O}_{4}$ | bend | $\sigma$ |
| 10 | $\mathrm{C}_{1} \mathrm{C}_{2} \mathrm{H}_{5}$ | bend | $\rho$ |
| 11 | $\mathrm{C}_{3} \mathrm{C}_{2} \mathrm{H}_{5}$ | bend | $\delta$ |
| 12 | $\mathrm{C}_{2} \mathrm{C}_{1} \mathrm{H}_{6}$ | bend | $\alpha_{1}$ |
| 13 | $\mathrm{C}_{2} \mathrm{C}_{1} \mathrm{~F}_{7}$ | bend | $\alpha_{2}$ |
| 14 | $\mathrm{C}_{2} \mathrm{C}_{1} \mathrm{~F}_{8}$ | bend | $\alpha_{3}$ |
| 15 | $\mathrm{H}_{6} \mathrm{C}_{1} \mathrm{~F}_{7}$ | bend | $\beta_{1}$ |
| 16 | $\mathrm{H}_{6} \mathrm{C}_{1} \mathrm{~F}_{8}$ | bend | $\beta_{2}$ |
| 17 | $\mathrm{F}_{7} \mathrm{C}_{1} \mathrm{~F}_{8}$ | bend | $\beta_{3}$ |
| 18 | $\mathrm{C}_{2}=\mathrm{C}_{3}=\mathrm{O}_{4}$ | wag | $\chi$ |
| 19 | $\mathrm{H}_{5} \mathrm{C}_{1} \mathrm{C}_{2} \mathrm{C}_{3}$ | wag | $\pi$ |
| 20 | $\begin{aligned} & \mathrm{H}_{6} \mathrm{C}_{1} \mathrm{C}_{2} \mathrm{C}_{3} \\ & \mathrm{~F}_{7} \mathrm{C}_{1} \mathrm{C}_{2} \mathrm{C}_{3} \end{aligned}$ | $\left(\mathrm{CHF}_{2}\right)$ torsion | $\tau$ |
|  | $\mathrm{F}_{8} \mathrm{C}_{1} \mathrm{C}_{2} \mathrm{C}_{3}$ |  |  |

$\left(\partial \sigma_{j} / \partial \Omega\right)$, which are proportional to the intensities, [57, 58] are given as:

$$
\begin{aligned}
& \partial \sigma_{j} / \partial \Omega=\left(2^{4} \pi^{4} / 45\right)\left(v_{\mathrm{o}}-v_{j}\right)^{4}\left(h / 8 \pi^{2} c v_{j}\right) S_{j}\left[\left(1-\rho_{j}\right) /\right. \\
&\left.\cdot\left(1+\rho_{j}\right)\right] /\left[1-\exp \left(-h c v_{j} / k_{\mathrm{B}} T\right)\right]
\end{aligned}
$$

Since we use only relative intensities, we calculated them as:
$I_{j}=\left(\partial \sigma_{\mathrm{j}} / \partial \Omega\right) /\left(\partial \sigma_{j m} / \partial \Omega\right)$
where $j m$ denotes the normal mode having the largest Raman cross-section. As laser wavelength we took $\lambda_{0}=514.5 \mathrm{~nm}\left(\mathrm{v}_{0}=1 / \lambda_{0}\right)$, which corresponds to an argon ion laser. We assumed the temperature to be $T=300 \mathrm{~K}$ in the case of 3,3-difluoromethyl ketene.

The line shapes were then calculated as Lorentzians ( $L$ ) with a width of $\Delta v=12 \mathrm{~cm}^{-1}$. Thus, the final spectrum is calculated as:

$$
\begin{aligned}
I(v) & =\sum_{j} I_{j} L\left(v-v_{j}\right) \\
L\left(v-v_{j}\right) & =(1 / \pi)(\Delta v / 2) /\left[\left(v-v_{j}\right)^{2}+(\Delta v / 2)^{2}\right] \\
\int_{-\infty}^{+\infty} L(v) \mathrm{d} v & =1
\end{aligned}
$$

where $j$ runs over all normal modes. For the plots we used a grid step size of $10 \mathrm{~cm}^{-1}$, but not when a spectral line appeared between two consecutive grid points. In this case we inserted 12 points with step size $0.5 \mathrm{~cm}^{-1}$ into this interval which include the exact location of the center of the line.

Table 7 Symmetry coordinates for trifluoromethyl ketene

| Species | Description | Symmetry coordinate ${ }^{\text {a }}$ |
| :---: | :---: | :---: |
| A' | C-H stretch | $\mathrm{S}_{1}=\mathrm{D}$ |
|  | $\mathrm{C}-\mathrm{C}$ stretch | $\mathrm{S}_{2}=\mathrm{R}$ |
|  | $\mathrm{C}=\mathrm{C}=\mathrm{O}$ symmetric stretch | $\mathrm{S}_{3}=\mathrm{P}+\mathrm{T}$ |
|  | $\mathrm{C}=\mathrm{C}=\mathrm{O}$ antisymmetric stretch | $\mathrm{S}_{4}=\mathrm{P}-\mathrm{T}$ |
|  | $\mathrm{CF}_{3}$ symmetric stretch | $\mathrm{S}_{5}=\mathrm{Q}_{1}+\mathrm{Q}_{2}+\mathrm{Q}_{3}$ |
|  | $\mathrm{CF}_{3}$ antisymmetric stretch | $\mathrm{S}_{6}=2 \mathrm{Q}_{1}-\mathrm{Q}_{2}-\mathrm{Q}_{3}$ |
|  | $\mathrm{C}=\mathrm{C}-\mathrm{H}$ in-plane bend | $\mathrm{S}_{7}=\rho-\delta$ |
|  | $\mathrm{C}=\mathrm{C}-\mathrm{C}$ in-plane bend | $\mathrm{S}_{8}=2 \gamma-\rho-\delta$ |
|  | $\mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bend | $\mathrm{S}_{9}=\sigma$ |
|  | $\mathrm{CF}_{3}$ symmetric deformation | $\mathrm{S}_{10}=\alpha_{1}+\alpha_{2}+\alpha_{3}-\beta_{1}-\beta_{2}-\beta_{3}$ |
|  | $\mathrm{CF}_{3}$ antisymmetric deformation | $\mathrm{S}_{11}=\beta_{1}+\beta_{2}-2 \beta_{3}$ |
|  | $\mathrm{CF}_{3}$ rock | $\mathrm{S}_{12}=2 \alpha_{1}-\alpha_{2}-\alpha_{3}$ |
| A" | $\mathrm{C}=\mathrm{C}-\mathrm{H}$ out-of-plane bend (wag) | $\mathrm{S}_{13}=\pi$ |
|  | $\mathrm{C}=\mathrm{C}=\mathrm{O}$ out-of-plane bend (wag) | $\mathrm{S}_{14}=\chi$ |
|  | $\mathrm{CF}_{3}$ antisymmetric stretch | $\mathrm{S}_{15}=\mathrm{Q}_{2}-\mathrm{Q}_{3}$ |
|  | $\mathrm{CF}_{3}$ antisymmetric deformation | $\mathrm{S}_{16}=\beta_{1}-\beta_{2}$ |
|  | $\mathrm{CF}_{3}$ rock | $\mathrm{S}_{17}=\alpha_{2}-\alpha_{3}$ |
|  | $\mathrm{CF}_{3}$ torsion | $\mathrm{S}_{18}=\tau$ |

a Not normalized

For the infrared spectrum we used the intensities as given by the DFT-B3LYP/6-311++G** calculations and converted them to relative transmittance. The Boltzmann distribution was then used to superimpose the spectra of the cis and the gauche conformers of difluoromethyl ketene. Accordingly, the total line intensity was calculated as a function of frequency. For $N+1$ different energy minima we have:
$I(v)=\sum_{l=0}^{N} g_{l} I_{l}(v) e^{-\Delta E_{l} / k_{R} T} / \sum_{l=0}^{N} g_{l} e^{-\Delta E_{l} / k_{B} T}$
where $I_{l}(v)$ are the line intensities for the respective minima, $\Delta E_{l}=E_{l}-E_{0}$, and $g_{l}$ is a degeneracy factor. The $g_{l}$ value for the trifluoro- and difluoromethyl ketene in the cis conformation is unity, whereas it is 2 for the nonplanar gauche conformer. Additionally, $E_{l}$ represents the total energy of conformer $l$ corrected with the corresponding zero-point energy. The label $l=0$ is given to the lowest energy conformer. For the temperature, $T=300 \mathrm{~K}$ was used. The calculated vibrational Raman and infrared


Fig. 3 Calculated vibrational IR spectrum of trifluoromethyl ketene at the DFT-B3LYP/6-311++G** level

Table 8 Symmetry coordinates for difluoromethyl ketene

| Species | Description | Symmetry coordinate ${ }^{\text {a }}$ |
| :---: | :---: | :---: |
| A' | C- $\mathrm{H}_{5}$ stretch | $\mathrm{S}_{1}=\mathrm{D}$ |
|  | $\mathrm{C}-\mathrm{C}$ stretch | $\mathrm{S}_{2}=\mathrm{R}$ |
|  | $\mathrm{C}=\mathrm{C}=\mathrm{O}$ symmetric stretch | $\mathrm{S}_{3}=\mathrm{P}+\mathrm{T}$ |
|  | $\mathrm{C}=\mathrm{C}=\mathrm{O}$ antisymmetric stretch | $\mathrm{S}_{4}=\mathrm{P}-\mathrm{T}$ |
|  | $\mathrm{CF}_{2}$ symmetric stretch | $\mathrm{S}_{5}=\mathrm{Q}_{2}+\mathrm{Q}_{3}$ |
|  | $\mathrm{C}-\mathrm{H}_{6}$ stretch | $\mathrm{S}_{6}=\mathrm{Q}_{1}$ |
|  | $\mathrm{C}=\mathrm{C}-\mathrm{H}$ in-plane bend | $\mathrm{S}_{7}=\rho-\delta$ |
|  | $\mathrm{C}=\mathrm{C}-\mathrm{C}$ in-plane bend | $\mathrm{S}_{8}=2 \gamma-\rho-\delta$ |
|  | $\mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bend | $\mathrm{S}_{9}=\sigma$ |
|  | $\mathrm{CH}_{6}$ in-plane bend | $S_{10}=\left[(6)^{1 / 2}-2\right] \beta_{3}-\left[(6)^{1 / 2}+2\right] \alpha_{1}+\alpha_{2}+\alpha_{3}+\beta_{1}+\beta_{2}$ |
|  | $\mathrm{CF}_{2}$ deformation | $\mathrm{S}_{11}=\left[(6)^{1 / 2}+2\right] \beta_{3}-\left[(6)^{1 / 2}-2\right] \alpha_{1}-\alpha_{2}-\alpha_{3}-\beta_{1}-\beta_{2}$ |
|  | $\mathrm{CF}_{2}$ rock | $\mathrm{S}_{12}=\alpha_{2}+\alpha_{3}-\beta_{1}-\beta_{2}$ |
| A" | $\mathrm{C}=\mathrm{C}-\mathrm{H}$ out-of-plane bend (wag) | $\mathrm{S}_{13}=\pi$ |
|  | $\mathrm{C}=\mathrm{C}=\mathrm{O}$ out-of-plane bend (wag) | $\mathrm{S}_{14}=\chi$ |
|  | $\mathrm{CF}_{2}$ antisymmetric stretch | $\mathrm{S}_{15}=\mathrm{Q}_{2}-\mathrm{Q}_{3}$ |
|  | $\mathrm{C}-\mathrm{H}_{6}$ out-of-plane bend | $\mathrm{S}_{16}=\alpha_{2}-\alpha_{3}+\beta_{1}-\beta_{2}$ |
|  | $\mathrm{CF}_{2}$ twist | $\mathrm{S}_{17}=\alpha_{2}-\alpha_{3}-\beta_{1}+\beta_{2}$ |
|  | $\mathrm{CHF}_{2}$ torsion | $\mathrm{S}_{18}=\tau$ |

Table 9 Calculated vibrational frequencies $\left(\mathrm{cm}^{-1}\right)$ at B3LYP/6-311++G** level for the cis conformer of trifluoromethyl ketene

| Symm. | No. | Frequency | IR intensity ${ }^{\text {a }}$ | Raman activity | Depol. ratio | Observed ${ }^{\text {b }}$ | PED ${ }^{\text {c }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| A' | $\mathrm{V}_{1}$ | 3215 | 26.4 | 77.6 | 0.2 | 3160 | ```100% C-H str. (S S ) 97% C=C=O antisymm. str. (S S ) 43% C=C=O symm. str. ( }\mp@subsup{\textrm{S}}{3}{}),25% C=C-H in-plane bend (S S ) 24% C-C str. (S S``` |
|  | $v_{2}$ | 2252 | 714.5 | 12.6 | 0.7 |  |  |
|  | $v_{3}$ | 1442 | 154.3 | 8.4 | 0.1 |  |  |
|  | $\mathrm{V}_{4}$ | 1235 | 256.3 | 0.4 | 0.7 | 1298 | $17 \% \mathrm{CF}_{3}$ symm. str. $\left(\mathrm{S}_{5}\right), 22 \% \mathrm{C}-\mathrm{C}$ str. $\left(\mathrm{S}_{2}\right), 17 \% \mathrm{CF}_{3}$ symm. def. $\left(\mathrm{S}_{10}\right)$, $17 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ in-plane bend $\left(\mathrm{S}_{7}\right), 14 \% \mathrm{CF}_{3}$ antisymm. str. $\left(\mathrm{S}_{6}\right)$ |
|  | $\mathrm{V}_{5}$ | 1136 | 265.2 | 6.5 | 0.3 | 1197 | $22 \% \mathrm{CF}_{3}$ antisymm. str. $\left(\mathrm{S}_{6}\right), 36 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ symm. str. $\left(\mathrm{S}_{3}\right)$, $17 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ in-plane bend $\left(\mathrm{S}_{7}\right), 11 \% \mathrm{CF}_{3}$ symm. str. $\left(\mathrm{S}_{5}\right)$ |
|  | $v_{6}$ | 1100 | 155.2 | 11.5 | 0.3 |  | $35 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ in-plane bend $\left(\mathrm{S}_{7}\right), 34 \% \mathrm{CF}_{3}$ antisymm. str. $\left(\mathrm{S}_{6}\right)$ |
|  | $v_{7}$ | 856 | 12.1 | 6.5 | 0.1 | 841 | $45 \% \mathrm{CF}_{3}$ symm. str. $\left(\mathrm{S}_{5}\right), 21 \% \mathrm{C}-\mathrm{C}$ str. $\left(\mathrm{S}_{2}\right)$, <br> $10 \% \mathrm{CF}_{3}$ antisymm. str. $\left(\mathrm{S}_{6}\right), 10 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bend $\left(\mathrm{S}_{9}\right)$ |
|  | $v_{8}$ | 709 | 4.7 | 5.3 | 0.2 | 707 | $32 \% \mathrm{CF}_{3}$ symm. def. $\left(\mathrm{S}_{10}\right), 23 \% \mathrm{CF}_{3}$ symm. str. $\left(\mathrm{S}_{5}\right)$, $18 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bend $\left(\mathrm{S}_{9}\right), 12 \% \mathrm{C}=\mathrm{C}-\mathrm{C}$ in-plane bend $\left(\mathrm{S}_{8}\right)$ |
|  | $v_{9}$ | 535 | 0.6 | 0.6 | 0.7 | 531 | $69 \% \mathrm{CF}_{3}$ antisymm. def. $\left(\mathrm{S}_{11}\right), 12 \% \mathrm{CF}_{3}$ antisymm. str. $\left(\mathrm{S}_{6}\right)$ $30 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bend $\left(\mathrm{S}_{9}\right), 33 \% \mathrm{CF}_{3}$ symm. def. $\left(\mathrm{S}_{10}\right)$, $27 \% \mathrm{C}-\mathrm{C}$ str. $\left(\mathrm{S}_{2}\right)$ |
|  | $v_{10}$ | 509 | 13.9 | 2.4 | 0.1 |  |  |
|  | $\mathrm{v}_{11}$ | 361 | 3.1 | 1.5 | 0.6 | 255 | $62 \% \mathrm{CF}_{3}$ rock $\left(\mathrm{S}_{12}\right), 17 \% \mathrm{CF}_{3}$ antisymm. def. $\left(\mathrm{S}_{11}\right)$, $11 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bend $\left(\mathrm{S}_{9}\right)$ |
|  | $v_{12}$ | 142 | 2.6 | 1.7 | 0.7 |  | $65 \% \mathrm{C}=\mathrm{C}-\mathrm{C}$ in-plane bend $\left(\mathrm{S}_{8}\right), 19 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bend $\left(\mathrm{S}_{9}\right)$, $16 \% \mathrm{CF}_{3}$ rock $\left(\mathrm{S}_{12}\right)$ |
| A" | $V_{13}$ | 1101 | 355.8 | 2.8 | 0.7 | 1194 | $77 \% \mathrm{CF}_{3}$ antisymm. str. $\left(\mathrm{S}_{15}\right), 13 \% \mathrm{CF}_{3}$ antisymm. def. $\left(\mathrm{S}_{16}\right)$ $23 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ wag $\left(\mathrm{S}_{13}\right), 31 \% \mathrm{CF}_{3}$ rock $\left(\mathrm{S}_{17}\right)$, <br> $15 \% \mathrm{CF}_{3}$ antisymm. str. $\left(\mathrm{S}_{15}\right), 14 \% \mathrm{CF}_{3}$ torsion $\left(\mathrm{S}_{18}\right)$ |
|  | $v_{14}$ | 630 | 28.0 | 0.6 | 0.7 |  |  |
|  | $v_{15}$ | 547 | 2.5 | 1.3 | 0.7 |  |  |
|  | $\nu_{16}$ | 508 | 9.4 | 1.3 | 0.7 | 529 d | $56 \% \mathrm{CF}_{3}$ antisymm. def. $\left(\mathrm{S}_{16}\right), 22 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ wag $\left(\mathrm{S}_{13}\right)$, $14 \% \mathrm{CF}_{3}$ torsion $\left(\mathrm{S}_{18}\right)$ |
|  | $v_{17}$ | 324 | 12.6 | 0.7 | 0.7 | 320 | $60 \% \mathrm{CF}_{3}$ rock $\left(\mathrm{S}_{17}\right), 25 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ wag $\left(\mathrm{S}_{13}\right)$, $15 \% \mathrm{CF}_{3}$ antisymm. def. $\left(\mathrm{S}_{16}\right)$ |
|  | $\mathrm{V}_{18}$ | 50 | 2.2 | 2.4 | 0.7 | 65 | $74 \% \mathrm{CF}_{3}$ torsion $\left(\mathrm{S}_{18}\right), 26 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ wag $\left(\mathrm{S}_{13}\right)$ |

${ }^{\text {a }}$ Infrared intensities and Raman activities are calculated in $\mathrm{Km} \mathrm{mol}^{-1}$ and $\AA^{4} \mathrm{amu}^{-1}$, respectively
${ }^{\mathrm{b}}$ Observed Raman frequencies $\left(\mathrm{cm}^{-1}\right)$ in the gas phase for trifluoroacetaldehyde (fluoral) are obtained from reference [60]


Fig. 4 Calculated vibrational Raman spectrum of trifluoromethyl ketene at the DFT-B3LYP/6-311++G** level
spectra of the cis-trifluoromethyl ketene are shown in Figs. 3 and 4, while the spectra of the difluoromethyl ketene mixture ( $65 \%$ cis and $35 \%$ gauche) at 300 K are shown in Figs. 5 and 6.
c PED values are obtained by using calculated frequencies at B3LYP level
${ }^{\mathrm{d}}$ IR frequency


Fig. 5 Calculated vibrational IR spectrum of difluoromethyl ketene at 300 K at the DFT-B3LYP/6-311++G** level

## Discussion

The interesting properties of organoisocyanate molecules $[33,34,35,36,37,38,39,40,41,42,43,44,45$, $46,47,48,49,50,51]$ recently prompted us to investigate the conformational and structural stabilities of

Table 10 Calculated vibrational frequencies $\left(\mathrm{cm}^{-1}\right)$ at B3LYP/6-311++G** level for the cis conformer of difluoromethyl ketene

| Symm. | No. | Frequency | IR intensitya | Raman activity | Depol. ratio | PED ${ }^{\text {b }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| A' | $v_{1}$ | 3196 | 19.9 | 76.5 | 0.2 | $100 \% \mathrm{C}-\mathrm{H}_{5}$ str. ( $\mathrm{S}_{1}$ ) |
|  | $v_{2}$ | 3100 | 24.8 | 73.5 | 0.2 | $100 \% \mathrm{C}-\mathrm{H}_{6}$ str. ( $\mathrm{S}_{6}$ ) |
|  | $v_{3}$ | 2234 | 752.7 | 13.1 | 0.7 | 97\% $\mathrm{C}=\mathrm{C}=\mathrm{O}$ antisymm. str. $\left(\mathrm{S}_{4}\right)$ |
|  | $\mathrm{v}_{4}$ | 1466 | 108.6 | 4.2 | 0.1 | $24 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ symm. str. $\left(\mathrm{S}_{3}\right), 34 \% \mathrm{C}-\mathrm{H}_{6}$ in-plane bend $\left(\mathrm{S}_{10}\right)$, ) <br> $15 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ in-plane bend $\left(\mathrm{S}_{7}\right), 14 \% \mathrm{C}-\mathrm{C}$ str. $\left(\mathrm{S}_{2}\right), 13 \% \mathrm{CF}_{2}$ rock $\left(\mathrm{S}_{12}\right.$ |
|  | $v_{5}$ | 1332 | 1.3 | 7.6 | 0.2 | $23 \% \mathrm{C}-\mathrm{H}_{6}$ in-plane bend $\left(\mathrm{S}_{10}\right), 31 \% \mathrm{CF}_{2}$ rock ( $\mathrm{S}_{12}$ ), <br> $29 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ symm. str. $\left(\mathrm{S}_{3}\right), 13 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ in-plane bend $\left(\mathrm{S}_{7}\right)$ |
|  | $v_{6}$ | 1143 | 7.8 | 7.5 | 0.6 | $55 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ in-plane bend $\left(\mathrm{S}_{7}\right), 17 \% \mathrm{C}-\mathrm{C}$ str. $\left(\mathrm{S}_{2}\right)$, $16 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ symm. str. $\left(\mathrm{S}_{3}\right)$ |
|  | $\mathrm{v}_{7}$ | 1087 | 229.1 | 4.0 | 0.1 | $40 \% \mathrm{C}-\mathrm{C}$ str. $\left(\mathrm{S}_{2}\right), 18 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ symm. str. $\left(\mathrm{S}_{3}\right), 15 \% \mathrm{CF}_{2}$ symm. str. $\left(\mathrm{S}_{5}\right)$ |
|  | $\mathrm{v}_{8}$ | 1052 | 141.7 | 14.3 | 0.2 | $60 \% \mathrm{CF}_{2}$ symm. str. ( $\mathrm{S}_{5}$ ) |
|  | $\mathrm{v}_{9}$ | 630 | 3.0 | 4.9 | 0.1 | $41 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bend $\left(\mathrm{S}_{9}\right), 17 \% \mathrm{C}=\mathrm{C}-\mathrm{C}$ in-plane bend $\left(\mathrm{S}_{8}\right)$, $13 \% \mathrm{C}-\mathrm{C}$ str. $\left(\mathrm{S}_{2}\right), 12 \% \mathrm{CF}_{2}$ symm. str. $\left(\mathrm{S}_{5}\right), 10 \% \mathrm{CF}_{2}$ def. ( $\mathrm{S}_{11}$ ) |
|  | $v_{10}$ | 565 | 3.4 | 0.8 | 0.6 | $26 \% \mathrm{CF}_{2}$ rock $\left(\mathrm{S}_{12}\right), 29 \% \mathrm{CF}_{2}$ def. $\left(\mathrm{S}_{11}\right), 18 \% \mathrm{C}=\mathrm{C}=0$ in-plane bend $\left(\mathrm{S}_{9}\right)$, $11 \% \mathrm{CF}_{2}$ symm. str. $\left(\mathrm{S}_{5}\right), 11 \% \mathrm{C}-\mathrm{H}_{6}$ in-plane bend ( $\mathrm{S}_{10}$ ) |
|  | $v_{11}$ | 426 | 22.2 | 2.2 | 0.3 | $47 \% \mathrm{CF}_{2}$ def. $\left(\mathrm{S}_{11}\right), 17 \% \mathrm{CF}_{2}$ rock $\left(\mathrm{S}_{12}\right), 15 \% \mathrm{C}-\mathrm{H}_{6}$ in-plane bend $\left(\mathrm{S}_{10}\right)$, $14 \% \mathrm{C}-\mathrm{C}$ str. ( $\mathrm{S}_{2}$ ) |
|  | $v_{12}$ | 173 | 8.7 | 2.5 | 0.7 | 66\% $\mathrm{C}=\mathrm{C}-\mathrm{C}$ in-plane bend ( $\mathrm{S}_{8}$ ), 28\% $\mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bend ( $\mathrm{S}_{9}$ ) |
| A" | $v_{13}$ | 1346 | 9.8 | 3.8 | 0.8 | $67 \% \mathrm{C}-\mathrm{H}_{6}$ wag ( $\mathrm{S}_{16}$ ), 33\% $\mathrm{CF}_{2}$ twist ( $\mathrm{S}_{17}$ ) |
|  | $\mathrm{v}_{14}$ | 1008 | 274.0 | 4.2 | 0.8 | 94\% $\mathrm{CF}_{2}$ antisymm. str. ( $\mathrm{S}_{15}$ ) |
|  | $v_{15}$ | 622 | 37.0 | 1.7 | 0.8 | $15 \% \mathrm{CF}_{2}$ twist $\left(\mathrm{S}_{17}\right), 33 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ wag $\left(\mathrm{S}_{13}\right)$, $28 \% \mathrm{CHF}_{2}$ torsion $\left(\mathrm{S}_{18}\right), 11 \% \mathrm{C}-\mathrm{H}_{6}$ wag $\left(\mathrm{S}_{16}\right)$ |
|  | $v_{16}$ | 548 | 0.1 | 0.6 | 0.8 | 94\% $\mathrm{C}=\mathrm{C}=\mathrm{O}$ wag ( $\mathrm{S}_{14}$ ) |
|  | $v_{17}$ | 346 | 14.5 | 1.2 | 0.8 | $49 \% \mathrm{CF}_{2}$ twist ( $\mathrm{S}_{17}$ ), 28\% C=C-H wag ( $\mathrm{S}_{13}$ ), 21\% $\mathrm{C}-\mathrm{H}_{6}$ wag ( $\mathrm{S}_{16}$ ) |
|  | $v_{18}$ | 63 | 4.1 | 2.3 | 0.8 | 65\% $\mathrm{CHF}_{2}$ torsion ( $\mathrm{S}_{18}$ ), 35\% $\mathrm{C}=\mathrm{C}-\mathrm{H}$ wag ( $\mathrm{S}_{13}$ ) |

${ }^{\mathrm{a}}$ Infrared intensities and Raman activities are calculated in $\mathrm{Km} \mathrm{mol}^{-1}$ and $\AA^{4} \mathrm{amu}^{-1}$ respectively
${ }^{\mathrm{b}}$ PED values are obtained by using calculated frequencies at B3LYP level

Table 11 Calculated vibrational frequencies $\left(\mathrm{cm}^{-1}\right)$ at B3LYP/6-311++G** level for the gauche conformer of difluoromethyl ketene

| Symm. | No. | Frequency | IR intensity ${ }^{\text {a }}$ | Raman activity | Depol. ratio | PED ${ }^{\text {b }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| A | $v_{1}$ | 3206 | 16.9 | 84.9 | 0.2 | $100 \% \mathrm{C}-\mathrm{H}_{5}$ str. ( $\mathrm{S}_{1}$ ) |
|  | $\mathrm{v}_{2}$ | 3090 | 34.3 | 120.7 | 0.2 | $67 \% \mathrm{C}-\mathrm{H}_{6}$ str. ( $\mathrm{S}_{6}$ ), 33\% $\mathrm{CF}_{2}$ symm. str. ( $\mathrm{S}_{5}$ ) |
|  | $v_{3}$ | 2239 | 753.6 | 12.2 | 0.7 | 97\% $\mathrm{C}=\mathrm{C}=\mathrm{O}$ antisymm. str. $\left(\mathrm{S}_{4}\right)$ |
|  | $v_{4}$ | 1432 | 40.6 | 8.1 | 0.1 | $44 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ symm. str. $\left(\mathrm{S}_{3}\right), 29 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ in-plane bend $\left(\mathrm{S}_{7}\right)$, $16 \% \mathrm{C}-\mathrm{C}$ str. $\left(\mathrm{S}_{2}\right)$ |
|  | $v_{5}$ | 1380 | 102.1 | 4.2 | 0.7 | $51 \% \mathrm{C}-\mathrm{H}_{6}$ in-plane bend $\left(\mathrm{S}_{10}\right), 19 \% \mathrm{C}-\mathrm{H}_{6}$ wag $\left(\mathrm{S}_{16}\right), 14 \% \mathrm{CF}_{2}$ rock $\left(\mathrm{S}_{12}\right)$ |
|  | $v_{6}$ | 1366 | 30.9 | 5.9 | 0.7 | $76 \% \mathrm{C}-\mathrm{H}_{6}$ wag ( $\mathrm{S}_{16}$ ), 14\% $\mathrm{C}-\mathrm{H}_{6}$ in-plane bend ( $\mathrm{S}_{10}$ ) |
|  | $v_{7}$ | 1146 | 48.0 | 12.5 | 0.3 | 49\% $\mathrm{C}=\mathrm{C}-\mathrm{H}$ in-plane bend ( $\mathrm{S}_{7}$ ), 13\% $\mathrm{C}=\mathrm{C}=\mathrm{O}$ symm. str. $\left(\mathrm{S}_{3}\right)$ |
|  | $\mathrm{v}_{8}$ | 1096 | 105.4 | 5.1 | 0.4 | $22 \% \mathrm{CF}_{2}$ symm. str. $\left(\mathrm{S}_{5}\right), 19 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ symm. str. $\left(\mathrm{S}_{3}\right)$, $14 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ in-plane bend $\left(\mathrm{S}_{7}\right), 10 \% \mathrm{C}-\mathrm{H}_{6}$ str. $\left(\mathrm{S}_{6}\right)$ |
|  | $v_{9}$ | 1014 | 291.1 | 7.3 | 0.4 | $71 \% \mathrm{CF}_{2}$ antisymm. str. ( $\mathrm{S}_{15}$ ), 11\% $\mathrm{CF}_{2}$ symm. str. ( $\mathrm{S}_{5}$ ) |
|  | $v_{10}$ | 915 | 32.7 | 8.0 | 0.04 | $42 \% \mathrm{C}-\mathrm{C}$ str. $\left(\mathrm{S}_{2}\right), 16 \% \mathrm{CF}_{2}$ rock $\left(\mathrm{S}_{12}\right), 12 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ symm. str. $\left(\mathrm{S}_{3}\right)$, $10 \% \mathrm{CF}_{2}$ symm. str. $\left(\mathrm{S}_{5}\right)$ |
|  | $v_{11}$ | 786 | 33.6 | 2.8 | 0.7 | $23 \% \mathrm{CF}_{2}$ rock $\left(\mathrm{S}_{12}\right), 21 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bend $\left(\mathrm{S}_{9}\right)$, <br> $20 \% \mathrm{C}=\mathrm{C}-\mathrm{C}$ in-plane bend $\left(\mathrm{S}_{8}\right)$ |
|  | $v_{12}$ | 598 | 29.2 | 3.2 | 0.2 | $11 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bend $\left(\mathrm{S}_{9}\right), 21 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ wag $\left(\mathrm{S}_{13}\right)$, $21 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ wag ( $\mathrm{S}_{14}$ ) |
|  | $v_{13}$ | 548 | 3.1 | 2.8 | 0.5 | $59 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ wag ( $\mathrm{S}_{14}$ ), 24\% $\mathrm{CF}_{2}$ def. $\left(\mathrm{S}_{11}\right)$ |
|  | $v_{14}$ | 519 | 6.0 | 0.8 | 0.7 | $23 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ wag $\left(\mathrm{S}_{13}\right), 29 \% \mathrm{CF}_{2}$ def. $\left(\mathrm{S}_{11}\right), 20 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ wag $\left(\mathrm{S}_{14}\right)$, $10 \% \mathrm{CHF}_{2}$ torsion ( $\mathrm{S}_{18}$ ), |
|  | $v_{15}$ | 429 | 10.2 | 0.7 | 0.6 | $30 \% \mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bend $\left(\mathrm{S}_{9}\right), 20 \% \mathrm{CF}_{2}$ twist $\left(\mathrm{S}_{17}\right), 18 \% \mathrm{CF}_{2}$ rock $\left(\mathrm{S}_{12}\right)$, $10 \% \mathrm{C}-\mathrm{H}_{6}$ in-plane bend $\left(\mathrm{S}_{10}\right)$ |
|  | $v_{16}$ | 361 | 17.8 | 1.3 | 0.6 | $49 \% \mathrm{CF}_{2}$ twist $\left(\mathrm{S}_{17}\right), 28 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ wag $\left(\mathrm{S}_{13}\right), 10 \% \mathrm{CF}_{2}$ rock $\left(\mathrm{S}_{12}\right)$ |
|  | $v_{17}$ | 143 | 1.1 | 1.6 | 0.7 | $65 \% \mathrm{C}=\mathrm{C}-\mathrm{C}$ in-plane bend ( $\mathrm{S}_{8}$ ), 20\% $\mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bend ( $\mathrm{S}_{9}$ ) |
|  | $\nu_{18}$ | 51 | 1.0 | 2.0 | 0.7 | $81 \% \mathrm{CHF}_{2}$ torsion $\left(\mathrm{S}_{18}\right), 18 \% \mathrm{C}=\mathrm{C}-\mathrm{H}$ wag $\left(\mathrm{S}_{13}\right)$ |

[^1]

Fig. 6 Calculated vibrational Raman spectrum of difluoromethyl ketene at 300 K at the DFT-B3LYP/6-311++G** level
isocyanatoacetaldehyde, chloromethyl isocyanate, difluoromethyl isocyanate and formyl isocyanate. [48, 49, 50, 51] The conformational behavior of these compounds depends greatly on the extent of interactions between the substituent and both the $\pi$-system of the isocyanate group and the lone pair on the nitrogen. In isocyanatoacetaldehyde the cis-cis conformation (isocyanate group eclipses the carbonyl group) was found to be the lowest energy form with the carbonyl oxygen being oriented away from the lone pair. [48] The potential surface governing internal rotation about the $\mathrm{C}-\mathrm{N}$ bond in chloromethyl isocyanate was calculated to be consistent with a single minimum corresponding to a structure having the chlorine atom gauche or near-cis to the -NCO moiety [49]. From a more recent study, difluoromethyl isocyanate was predicted to exist only in the trans conformation, again with minimum repulsive interaction between the fluorine atoms and the electron lone pair on the nitrogen. [50]

In the case of formyl isocyanate, the molecule was predicted to have a cis $\Leftrightarrow$ trans conformational equilibrium. [51] The molecule was found to have the cis (carbonyl oxygen eclipses the isocyanato group) conformer as the lower energy form, with the carbonyl oxygen being away from the electron lone pair. A similar conformational equilibrium was predicted for formyl ketene with a relatively higher rotational barrier. [9] The conjugation stabilization force in such systems stabilizes the planar forms and considerably restricts the internal rotation of the -CHO group.

The electronic structure of trifluoromethyl ketene fascinated us. On first sight, we expected that the trans conformer (a fluorine atom eclipses the ketenic hydrogen) should be more stable than the cis form (a fluorine atom eclipses the $\mathrm{C}=\mathrm{C}=\mathrm{O}$ moiety). However, closer inspection shows that actually the reverse was true. As shown in Table 12, the negatively charged $\mathrm{F}_{6}$ atom in the cis conformer is calculated to be quite close to the positive $\mathrm{C}_{3}$ atom. Moreover, in both conformers the repulsion between the negative $\mathrm{F}_{6}$ and $\mathrm{O}_{4}$ atoms is rather small due to the large distances between them. Thus, the $\mathrm{C}_{3}-\mathrm{F}_{6}$ and


Fig. 7a-e Orbital interactions in trifluoromethyl ketene. a Destabilizing interaction between two filled orbitals. b Stabilizing interaction between a filled and an unfilled orbital. c MO interaction of fluorine lone pairs (and C-F $\sigma$ orbitals) with the $\pi^{*}$ orbital of the $\mathrm{C}-\mathrm{C} \pi$-bond in the trans conformation of trifluoromethyl ketene. d MO interaction of fluorine lone pairs (and C-F $\sigma$ orbitals) with the $\pi^{*}$ orbital of the $\mathrm{C}-\mathrm{C} \pi$-bond in the cis conformation of trifluoromethyl ketene. e MO interaction of fluorine lone pairs (and C-F $\sigma$ orbitals) with the $\pi$ orbital of the $\mathrm{C}-\mathrm{C} \pi$-bond in the trans conformation of trifluoromethyl ketene
$\mathrm{C}_{3}-\mathrm{F}_{7}\left(\mathrm{~F}_{8}\right)$ electrostatic attractions dominate the $\mathrm{F}-\mathrm{O}_{4}$ repulsion in both cases. The corresponding stabilization is larger in the cis conformer than in the trans one, because the absolute values of the charges are larger and the C-F distance is smaller in the cis form. However, this simple electrostatic argument explains only why the cis is more stable than the trans form, not why cis is a minimum but trans a maximum in the potential curve (Fig. 2). A closer look at the MO interactions in the system provides an explanation for this result.

As is well known, the interactions between two filled orbitals destabilize (Fig. 7a), while those between a

Table 12 Total atomic charges and selected non-bonded distances calculated at B3LYP/6-311++G** level for cis and trans conformations of trifluoromethyl ketene and difluoromethyl ketene

| Atom | cis | trans ${ }^{\text {a }}$ | Atom | cis | trans |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Trifluoromethyl ketene |  |  | Difluoromethyl ketene |  |  |
| $\mathrm{C}_{1}$ | 0.098 | 0.125 | $\mathrm{C}_{1}$ | -0.261 | -0.210 |
| $\mathrm{C}_{2}$ | 0.029 | 0.048 | $\mathrm{C}_{2}$ | 0.141 | 0.065 |
| $\mathrm{C}_{3}$ | 0.189 | 0.135 | $\mathrm{C}_{3}$ | 0.102 | 0.193 |
| $\mathrm{O}_{4}$ | -0.179 | -0.166 | $\mathrm{O}_{4}$ | -0.181 | -0.175 |
| $\mathrm{H}_{5}$ | 0.220 | 0.218 | $\mathrm{H}_{5}$ | 0.225 | 0.173 |
| $\mathrm{F}_{6}$ | -0.164 | -0.137 | $\mathrm{H}_{6}$ | 0.224 | 0.220 |
| $\mathrm{F}_{7}$ | -0.099 | -0.112 | $\mathrm{F}_{7}$ | -0.125 | -0.133 |
| $\mathrm{F}_{8}$ | -0.099 | -0.112 | $\mathrm{F}_{8}$ | -0.125 | -0.133 |
| Non-bonded distance ( $\AA$ ) |  |  | Non-bonded distance ( $\AA$ ) |  |  |
| $\mathrm{F}_{6}-\mathrm{C}_{3}$ | 2.678 |  | $\mathrm{H}_{6}-\mathrm{C}_{3}$ | 2.624 |  |
| $\mathrm{F}_{6}-\mathrm{O}_{4}$ | 3.403 |  | $\mathrm{H}_{6}-\mathrm{O}_{4}$ | 3.416 |  |
| $\mathrm{F}_{8}-\mathrm{C}_{3}$ |  | 2.976 | $\mathrm{F}_{8}-\mathrm{C}_{3}$ |  | 2.963 |
| $\mathrm{F}_{8}-\mathrm{C}_{3}$ |  | 3.818 | $\mathrm{F}_{8}-\mathrm{O}_{4}$ |  | 3.793 |

${ }^{a}$ Transition state
filled and an unfilled orbital stabilize (Fig. 7b) a system. Candidates for such interactions in our system are the lone pairs at fluorine, the $\sigma, \sigma^{*}$ orbitals of the C-F bond and the $\pi, \pi^{*}$ orbitals of the $\mathrm{C}=\mathrm{C} \pi$ bond. Being flanked by a $\mathrm{CF}_{3}$ group at one end and a $\mathrm{C}=\mathrm{O}$ bond at the other, the $\mathrm{C}=\mathrm{C}$ bond can be considered as quite electron deficient, making it an ideal system to interact with F lone pairs or the $\sigma$ orbitals of the C-F bond. However, as Fig. 7c shows, in the trans conformer neither of them can efficiently overlap with the $\pi^{*}$ orbital due to geometry, while in the cis form $\mathrm{F}_{6}$ lies in the nodal plane of the $\pi^{*}$ orbital (Fig. 7d). Thus, there are no efficient stabilizing orbital interactions and we have to concentrate on the destabilizing ones. In the case of the destabilizing interactions there is a significant difference between the cis and the trans form. Again in the cis conformer the lone pair at $\mathrm{F}_{6}$ and the $\mathrm{C}-\mathrm{F}_{6} \sigma$ orbital are exactly in the nodal plane of the $\pi$ orbital and are symmetric with respect to this plane, while $\pi$ is antisymmetric. Therefore, in the cis form there are no efficient destabilizing orbital interactions present as compared to that in the trans conformer. Furthermore, any arbitrary rotation away from the fully symmetric cis form destroys the symmetry of the C-F $\sigma$ orbital or the $\mathrm{F}_{6}$ lone pair, increasing destabilizing interactions. Hence, cis must be a true energy minimum.

On the contrary, in the trans form, as shown in Fig. 7e, there are destabilizing interactions present, and thus the trans conformer has to be thermodynamically less preferable than the cis form. Additionally, Fig. 7e indicates that any rotation out of the fully symmetric trans form would reduce these destabilizing effects, and thus the trans form is a maximum in the potential curve. Given the fact that the distances between the atoms indicate comparatively small effects due to the discussed orbital interactions, one would expect to find small barriers as predicted in our calculations. The difference in MO interactions between the two forms we have just viewed above surely dominates any electrostatic attractive inter-
actions between the fluorine and the ketenic hydrogen atom because the latter are of roughly the same strength in both conformers.

An electron diffraction study for trifluoroacetaldehyde concluded that the cis conformation (fluorine eclipses the aldehyde group) was the preferred conformation for the system. [55] The results obtained from the electron diffraction experiment were reasonably consistent with the microwave study. [59] As a result this experimental conclusion supports our prediction for the stability of trifluoromethyl ketene.

In the case of the difluoro derivative, the steric effects play a competitive role with the MO orbital interactions in determining the conformational stability in difluoromethyl ketene. The question of which bond, $\mathrm{C}-\mathrm{H}_{6}$ or $\mathrm{C}-\mathrm{F}$, eclipses the ketenic hydrogen $\left(\mathrm{C}-\mathrm{H}_{5}\right)$ or the CCO group is also important, because the different effects of steric hindrance might outweigh those due to MO interactions. In difluoromethyl ketene the low minima are expected at such angles, for which either $\mathrm{CH}_{6}$ or CF eclipses the ketenic CCO group and the $\mathrm{C}-\mathrm{H}_{5}$ bond is staggered relative to the other two bonds of the difluoromethyl $\left(\mathrm{CHF}_{2}\right)$ rotor.

In the lowest energy cis conformation of difluoromethyl ketene, the small $\mathrm{H}_{6}$ atom opposes the CCO group while the more bulky $\mathrm{C}-\mathrm{F}$ bonds eclipse the $\mathrm{C}-\mathrm{H}_{5}$ bond. The second, lower energy gauche conformer with an HCCC dihedral angle of $131^{\circ}$ has the $\mathrm{C}-\mathrm{F}_{8}$ eclipsing the CCO moiety. It is higher in energy than the cis conformer, because here the bulky $\mathrm{F}_{8}$ opposes the CCO group. However, although in both forms the ketenic CCO group is in an eclipsed position, the two are the lower energy forms, because they are the only forms where no destabilizing interactions between the fluorine lone pairs and the $\pi$ orbitals of the $\mathrm{C}=\mathrm{C}$ double bond occur. In the cis case no lone pair points to the $\mathrm{C}=\mathrm{C}$ double bond, while in the higher gauche ( HCCC is $120^{\circ}$ ) case the $\mathrm{F}_{8}$ lone pair lies in the nodal plane of the $\mathrm{C}=\mathrm{C}$ bond (see Fig. 7d).

The other gauche conformer ( HCCC is $60^{\circ}$ ) and the trans structure both suffer from the aforementioned destabilizing interactions. The trans form has two such interactions, while the high energy gauche structure has only one. Since these interactions are rather small in difluoromethyl ketene, not the trans form is higher in energy, but the gauche structure ( HCCC of $60^{\circ}$ ). The reason for this is simply that the bulky $\mathrm{C}-\mathrm{F}_{7}$ bond eclipses $\mathrm{C}-\mathrm{H}_{5}$ in the case of the high energy gauche form, while in the trans case only $\mathrm{C}-\mathrm{H}_{6}$ is eclipsing $\mathrm{C}-\mathrm{H}_{5}$. While MO interactions would give the gauche-60 structure the lower energy, their effects are outweighed by the larger steric hindrance in the gauche form, making it higher in energy than the trans structure. Since in both gauche-60 and trans, any rotation out of their respective positions would lower the destabilizing effects and finally drive them into one of the minima, the two structures are maxima, with gauche-60 being higher in energy than trans.

We calculated the vibrational frequencies and derived the PED values among symmetry coordinates in the low
energy conformers of the two molecules. Some of the vibrations were clearly pure, and their assignments were easily made possible only based on the PED values as shown in Tables 9, 10 and 11, and others were predicted to be mixed with neighboring modes.

The $\mathrm{C}-\mathrm{H}$ stretching modes of the two molecules were calculated to have the highest Raman activities. Only one $\mathrm{C}-\mathrm{H}$ stretching mode is in the Raman spectrum of trifluoromethyl ketene and was calculated to be at $3,215 \mathrm{~cm}^{-1}$. While in difluoromethyl ketene Raman spectrum there are two $\mathrm{C}-\mathrm{H}$ stretches that were calculated at $3,196 \mathrm{~cm}^{-1}$ for the ketenic $\mathrm{C}-\mathrm{H}$ stretch and at $3,100 \mathrm{~cm}^{-1}$ for the other $\mathrm{C}-\mathrm{H}$ stretch for the cis conformer.

The highest calculated IR intensities in the tri- and difluoromethyl ketene IR spectra were assigned to the $-\mathrm{C}=\mathrm{C}=\mathrm{O}$ antisymmetric stretching vibrations with PED values of $97 \%$. Furthermore, the $\mathrm{C}=\mathrm{C}-\mathrm{C}$ in-plane bending modes of the cis conformers of both molecules were calculated at $142 \mathrm{~cm}^{-1}$ for the trifluoro derivative and $173 \mathrm{~cm}^{-1}$ for the difluoro one. The $-\mathrm{C}=\mathrm{C}=\mathrm{O}$ symmetric stretching mode in the trifluoro system was of a clear assignment and was calculated at $1,442 \mathrm{~cm}^{-1}$ with $43 \%$ PED. In the case of the difluoro-derivative, the assignment of this $-\mathrm{C}=\mathrm{C}=\mathrm{O}$ mode was not straightforward based on the calculated PED and could be assigned to either $v_{4}$ or $v_{5}$. However, the $-\mathrm{C}=\mathrm{C}=\mathrm{O}$ symmetric stretching mode could be assigned to the $v_{4}$ with the higher IR intensity as one might expect for such a vibration. Then, the assignment of $\nu_{5}$ could be predicted from the PED values in Table 9 to the $\mathrm{C}-\mathrm{H}_{6}$ in-plane bending.

Several vibrational modes in trifluoromethyl ketene were calculated to be mixed with other modes. The assignments of such modes could be clarified with the help of the observed frequencies obtained from the gas phase Raman spectrum of trifluoroacetaldehyde (fluoral). [60] For example, $v_{4}$ in our trifluoro system could be referred to either the $\mathrm{C}-\mathrm{C}$ stretching or the $\mathrm{CF}_{3}$ symmetric stretching. In fluoral, the experimental frequency assigned to the $\mathrm{CF}_{3}$ symmetric stretching mode was $1,298 \mathrm{~cm}^{-1}$. Therefore, $v_{4}$ in trifluoromethyl ketene could be assigned to $\mathrm{CF}_{3}$ symmetric stretching. However, the $\mathrm{C}-\mathrm{C}$ stretching vibration can be assigned with confidence to the $856 \mathrm{~cm}^{-1}$ frequency since the corresponding vibration in fluoral was observed at $841 \mathrm{~cm}^{-1}$. [60] Similarly, the $\mathrm{CF}_{3}$ antisymmetric stretching mode and the $\mathrm{C}=\mathrm{C}-\mathrm{H}$ in-plane bending mode were calculated to couple with each other in $\mathrm{v}_{5}$ and $\mathrm{v}_{6}$. Experimentally, the $\mathrm{CF}_{3}$ antisymmetric stretching mode was found at $1,197 \mathrm{~cm}^{-1}$, which is closer to the frequency of $v_{5}$ in trifluoromethyl ketene (Table 9). Thus, $v_{5}$ and $v_{6}$ can be assigned to the antisymmetric $\mathrm{CF}_{3}$ stretching and the inplane $\mathrm{C}=\mathrm{C}-\mathrm{H}$ bending modes, respectively.

The assignment of some of the $\mathrm{CF}_{3}$ vibrational modes below $1,000 \mathrm{~cm}^{-1}$ was obvious from derived PED values. The calculated wavenumbers at 535 and $361 \mathrm{~cm}^{-1}$ were referred to the $\mathrm{CF}_{3}$ antisymmetric deformation and the $\mathrm{CF}_{3}$ rock, respectively. The $\mathrm{CF}_{3}$ symmetric deformation in fluoral was observed at $707 \mathrm{~cm}^{-1}$; hence the $\mathrm{CF}_{3}$ symmetric deformation in trifluoromethyl ketene could
be assigned to $v_{8}$. This leaves us with the assignment of $v_{10}$ to the $-\mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bending vibration.

The calculated PEDs for difluoromethyl ketene in the cis conformation provided excellent help for the assignments of most of the skeletal vibrations (see Table 10). The frequencies $v_{6}, v_{7}, v_{8}, v_{9}$ and $v_{11}$ can with little doubt be assigned to the $\mathrm{C}=\mathrm{C}-\mathrm{H}$ in-plane bending ( $55 \%$ PED ), the $\mathrm{C}-\mathrm{C}$ stretching ( $40 \% \mathrm{PED}$ ), the $\mathrm{CF}_{2}$ symmetric stretching ( $60 \% \mathrm{PED}$ ), the $\mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bending ( $41 \% \mathrm{PED}$ ) and the $\mathrm{CF}_{2}$ deformation ( $47 \% \mathrm{PED}$ ), respectively. This should then assign $v_{10}$ at $565 \mathrm{~cm}^{-1}$ to the $\mathrm{CF}_{2}$ rocking mode. The remaining frequency at $1,332 \mathrm{~cm}^{-1}$ in the cis conformer of the difluoro system $\left(V_{5}\right)$ can be referred with confidence to the $\mathrm{C}-\mathrm{H}_{6}$ in-plane bending.

Many vibrational modes in the gauche form of the difluoro derivative, however, were difficult to assign based only on the calculated PEDs, especially for the region below $1,000 \mathrm{~cm}^{-1}$, at which a large degree of coupling of many bending modes was found (Table 11). From the experimental IR frequencies of chlorocarbonyl ketene obtained in the xenon matrix, [19] the $\mathrm{C}=\mathrm{C}-\mathrm{H}$ wag was observed at $531 \mathrm{~cm}^{-1}$ for the cis and $556 \mathrm{~cm}^{-1}$ for the trans conformers. The observed modes were weak in the IR spectrum. This information leads to the assignment of the $\mathrm{C}=\mathrm{C}-\mathrm{H}$ wag of the gauche difluoromethyl ketene to the depolarized wavenumber at $519 \mathrm{~cm}^{-1}$. Therefore, from the calculated PEDs, the $\mathrm{C}=\mathrm{C}=\mathrm{O}$ in-plane bending was assigned to $v_{12}$. Then, the $\mathrm{CF}_{2}$ rock and the $\mathrm{CF}_{2}$ twist could be assigned to $v_{11}$ and $v_{15}$, respectively. This approximate assignment leads one to assign $v_{8}$ at $1,096 \mathrm{~cm}^{-1}$ to the symmetric $\mathrm{CF}_{2}$ stretching.

Many of the bending modes of A" symmetry in the tri- and the difluoromethyl ketene spectra were calculated to have a very low degree of mixing and their assignments are shown in Tables 9 and 10. Among these modes only the $\mathrm{C}=\mathrm{C}-\mathrm{H}$ wag was of a low PED value in trifluoromethyl ketene. The $\mathrm{C}=\mathrm{C}-\mathrm{H}$ waging mode was predicted to be at 630 and $622 \mathrm{~cm}^{-1}$ in the triflouro and difluoro systems, respectively. The lowest vibrational mode in the spectra of the two molecules was the asymmetric torsion. The $\mathrm{CF}_{3}$ asymmetric torsion in the spectrum of trifluoromethyl ketene was calculated to be at $50 \mathrm{~cm}^{-1}$, and the $\mathrm{CHF}_{2}$ asymmetric torsion in the difluoro spectrum was predicted to be at $63 \mathrm{~cm}^{-1}$ in the cis form and $51 \mathrm{~cm}^{-1}$ in the gauche form. The asymmetric torsion modes were predicted to have a relatively small degree of contribution from the $\mathrm{C}=\mathrm{C}-\mathrm{H}$ waging mode as shown in Tables 9, 10 and 11.

Finally, we were able to provide reasonable vibrational assignments for the normal modes of the stable conformers of tri- and difluoromethyl ketene based on normal coordinate analysis. Good agreements were noticed when the calculated frequencies of trifluoromethyl ketene were compared to the experimental ones of fluoral, especially for the vibrational modes associated to the $\mathrm{CF}_{3}$ moiety.

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[^1]:    ${ }^{\text {a }}$ Infrared intensities and Raman activities are calculated in $\mathrm{Km} \mathrm{mol}^{-1}$ and $\AA^{4} \mathrm{amu}^{-1}$, respectively
    ${ }^{\mathrm{b}}$ PED values are obtained by using calculated frequencies at B3LYP level

